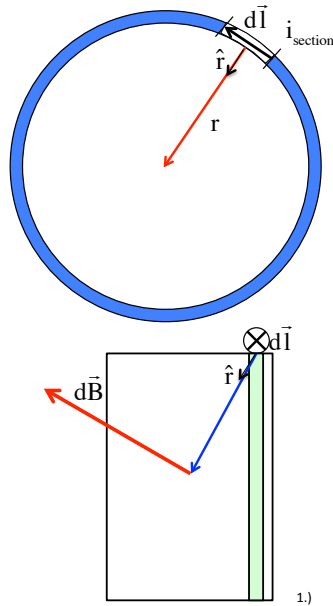


Problem 30.4

Define a strip of the induced current as i_{section} . From the side, a tiny section "dl" of that current would set up a magnetic field at the center of the cylinder defined by Biot Savart as:

$$dB = \frac{\mu_0 i_{\text{section}}}{4\pi} \frac{d\vec{l} \times \hat{r}}{r^2}$$

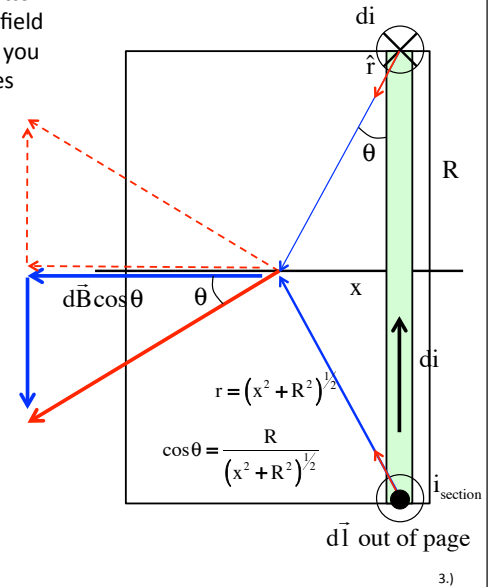
How this gets tricky can be seen by looking at the cylinder from the side. If the section associated with $d\vec{l}$ is right at the top of the page going into the page for the swath of current viewed, the right-thumb rule show that the direction of $d\vec{B}$ is seen not directed down the axis.



Look at the current section at the bottom of the swath. The vertical magnetic field components will add to zero leaving you with differential magnetic field pieces in the x-direction only. That means that the net field will be along the central axis of the hoop.

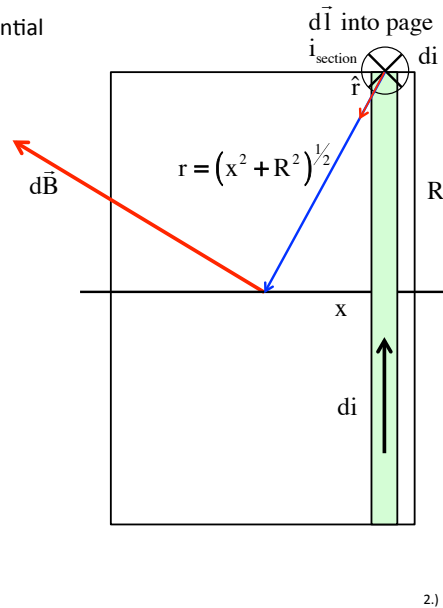
For our differential piece, we can write:

$$\begin{aligned} dB_x &= dB \cos \theta \\ &= \frac{\mu_0 (i_{\text{section}})(dl)}{4\pi(x^2 + R^2)} \cos \theta \\ &= \frac{\mu_0 (i_{\text{section}})(dl)}{4\pi(x^2 + R^2)} \frac{R}{(x^2 + R^2)^{1/2}} \\ &= \frac{\mu_0 R (i_{\text{section}})(dl)}{4\pi(x^2 + R^2)^{3/2}} \end{aligned}$$



According to Biot-Savart, that differential magnetic field vector will equal:

$$\begin{aligned} dB &= \frac{\mu_0 (i_{\text{sect}})}{4\pi r^2} d\vec{l} \times \hat{r} \\ &= \frac{\mu_0 (i_{\text{sect}})}{4\pi r^2} dl \sin 90^\circ \\ &= \frac{\mu_0 (i_{\text{sect}})}{4\pi r^2} dl \\ &= \frac{\mu_0 (i_{\text{sect}}) dl}{4\pi [(x^2 + R^2)^{1/2}]^2} \\ &= \frac{\mu_0 (i_{\text{sect}}) dl}{4\pi (x^2 + R^2)} \end{aligned}$$



If the width of the cylinder is L, the current density in **amps per meter** is such that

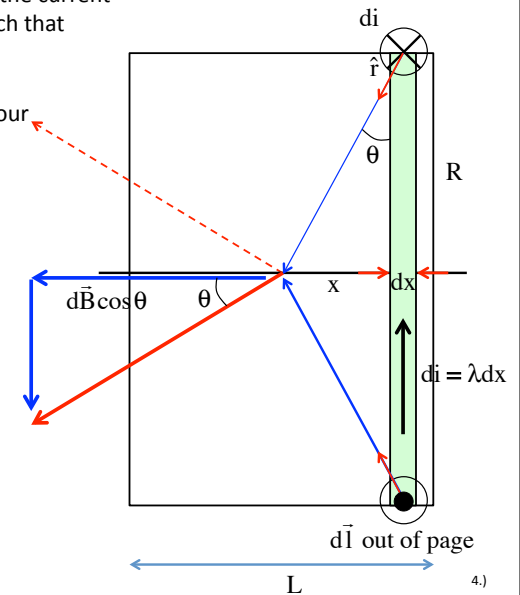
$$i_{\text{total}} = \lambda L \Rightarrow \lambda = i_{\text{total}} / L$$

With that, the current through our swath of width "dx" becomes:

$$i_{\text{section}} = \lambda dx$$

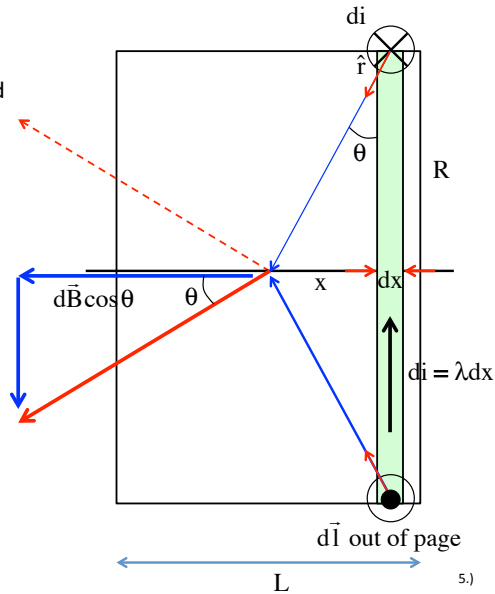
and

$$\begin{aligned} dB_x &= \frac{\mu_0 R (i_{\text{section}}) dl}{4\pi (x^2 + R^2)^{3/2}} \\ &= \frac{\mu_0 R (\lambda dx) dl}{4\pi (x^2 + R^2)^{3/2}} \end{aligned}$$



To get the total differential magnetic field due to the entire swath, we integrate over "dl" and get:

$$\begin{aligned} dB_{x,\text{total}} &= \frac{\mu_0 R (\lambda dx)}{4\pi(x^2 + R^2)^{3/2}} \oint dl \\ &= \frac{\mu_0 R (\lambda dx)}{4\pi(x^2 + R^2)^{3/2}} (2\pi R) \\ &= \frac{\mu_0 R^2 (\lambda dx)}{2(x^2 + R^2)^{3/2}} \end{aligned}$$

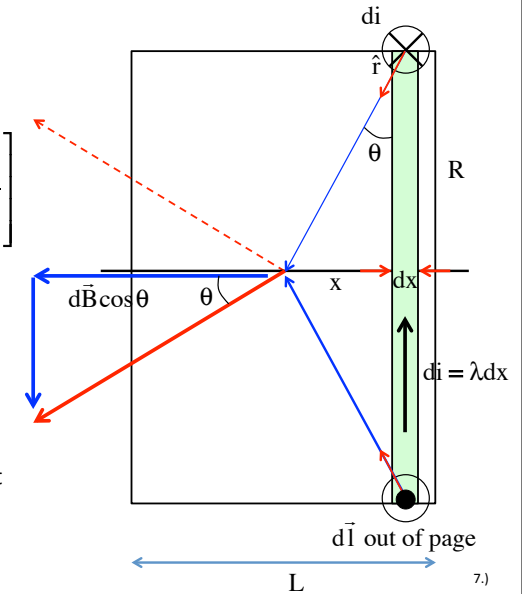


We get:

$$\begin{aligned} &= \mu_0 R^2 (\lambda) \left[\frac{x}{R^2(x^2 + R^2)^{1/2}} \right]_{x=0}^{L/2} \\ &= \mu_0 R^2 \left(\frac{i_{\text{total}}}{L} \right) \left[\frac{L/2}{R^2 \left((L/2)^2 + R^2 \right)^{1/2}} \right] \\ &= \frac{\mu_0 i_{\text{total}}}{2} \left[\frac{1}{\left((L/2)^2 + R^2 \right)^{1/2}} \right] \end{aligned}$$

(And kindly note that if L goes to zero, as would be the case if the cylinder was a hoop, we get the expected):

$$B = \frac{\mu_0 i_{\text{total}}}{2R}$$

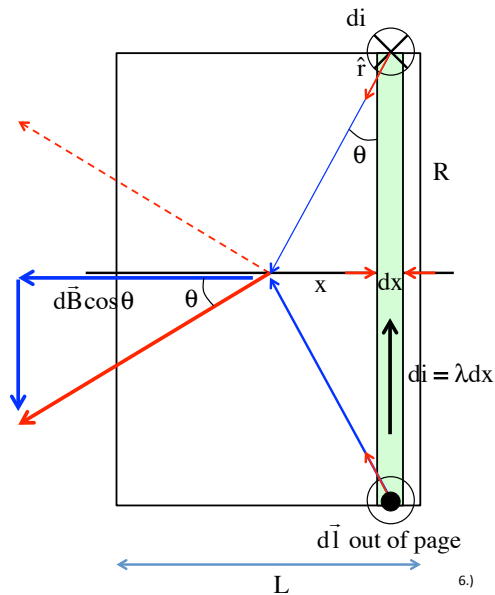


To get B due to all the differential swaths, we have to integrate over dx, or:

$$\begin{aligned} B &= 2 \int_{x=0}^{L/2} dB_{x,\text{total}} \\ &= 2 \left(\int_{x=0}^{L/2} \frac{\mu_0 R^2 (\lambda dx)}{2(x^2 + R^2)^{3/2}} \right) \\ &= (2) \left(\frac{\mu_0 R^2 (\lambda)}{2} \right) \int_{x=0}^{L/2} \frac{dx}{(x^2 + R^2)^{3/2}} \\ &= \mu_0 R^2 (\lambda) \left[\frac{x}{R^2(x^2 + R^2)^{1/2}} \right]_{x=0}^{L/2} \end{aligned}$$

Substituting in

$$\lambda = \frac{i_{\text{total}}}{L}$$



In any case, with our derived expression, the magnetic field at the center of the cylinder becomes:

$$\begin{aligned} B &= \frac{\mu_0 i}{2R} \\ &= \frac{\mu_0 \left(\frac{ev}{2\pi R} \right)}{2R} \\ &= \frac{\mu_0 ev}{4\pi R^2} \\ &= \frac{(4\pi \times 10^{-7} \text{ T} \cdot \text{m/A})(1.6 \times 10^{-19} \text{ C})(2.19 \times 10^6 \text{ m/s})}{4\pi (5.29 \times 10^{-11} \text{ m})^2} \\ &= 12.5 \text{ T} \end{aligned}$$

